

1 **Topographic shelf waves control seasonal melting near**
2 **Antarctic ice shelf grounding lines**

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10 **Key Points:**

- 11 • Variation in ice shelf basal melt rates is observed over a broad range of time scales,
12 from tidal to seasonal.
- 13 • Topographic shelf waves dominate the observed temporal melt rate variability in
14 the study region.
- 15 • Sea ice concentration and tidal currents modulate the magnitude and temporal
16 variation of melt rates.

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Abstract

The buttressing potential of ice shelves is modulated by changes in sub-shelf melting, in response to changing ocean conditions. We analyse the temporal variability in sub-shelf melting using an autonomous phase-sensitive radio-echo sounder (ApRES) near the grounding line of the Roi Baudouin Ice Shelf in East Antarctica. When combined with additional oceanographic evidence of seasonal variations in the stratification and the amplification of diurnal tides around the shelf break topography (Gunnerus Bank), the results suggest an intricate mechanism in which topographic waves control the seasonal melt rate variability near the grounding line. This mechanism has not been considered before, and has the potential to enhance local melt rates without advecting different water masses. As topographic waves seem to strengthen in a stratified ocean, the freshening of Antarctic surface water predicted by observations and models is likely to increase future basal melting in this area.

Plain Language Summary

Ice shelves (or the floating parts of the Antarctic ice sheet) lose primarily mass through melting at their bottom in contact with the ocean. This thins them and makes them more vulnerable to potential collapse. To understand the processes governing such thinning, direct and long-time measurements are essential. Here we report on the first high-resolution time series of direct melt measurements on the Roi Baudouin Ice Shelf in Dronning Maud Land during 2016. We find that sub-shelf melt varies on both seasonal and daily time scales. Temporal variations stem from topographical ocean waves that originate on the continental shelf and transfer ocean properties without time delay within the ice shelf cavity. Therefore, seasonal variations highly depend on the presence/absence of sea ice in front of the ice shelf, which impact the strength of topographical waves. This mechanism is highly efficient at increasing the ice-ocean exchanges and may explain regional differences in ice-shelf melt.

1 Introduction

More than 75% of Antarctic continental ice discharges through ice shelves [*Bindshadler et al.*, 2011] that buttress the grounded ice [*Dupont and Alley*, 2005]. A reduction in ice shelf buttressing may lead to grounded ice-flow acceleration [*Rack and Rott*, 2004; *Reese et al.*, 2018; *Schannwell et al.*, 2018]. Numerical ice-sheet models suggest that

48 ice-shelf thinning due to melting from their contact with warmer ocean waters could desta-
49 bilise the ice sheet [Favier *et al.*, 2014; Schannwell *et al.*, 2018], indicating a crucial de-
50 pendence of the ice sheet on processes at the ice-ocean interface. Here, we present new
51 and highly resolving time series of melt rates observed on the Roi Baudouin Ice Shelf (RBIS),
52 which is part of the belt of slowly melting [Rignot *et al.*, 2013] smaller ice shelves situ-
53 ated over the narrow continental shelf along the coast of Dronning Maud Land, East Antarc-
54 tica (Fig. 1). The observations are taken near the grounding line where the bed lowers
55 toward the interior of the ice sheet [Callens *et al.*, 2014; Fretwell *et al.*, 2013] and our
56 analysis explores the oceanographic mechanisms that control basal melting in this re-
57 gion that is potentially susceptible to marine ice sheet instability [Favier *et al.*, 2016].

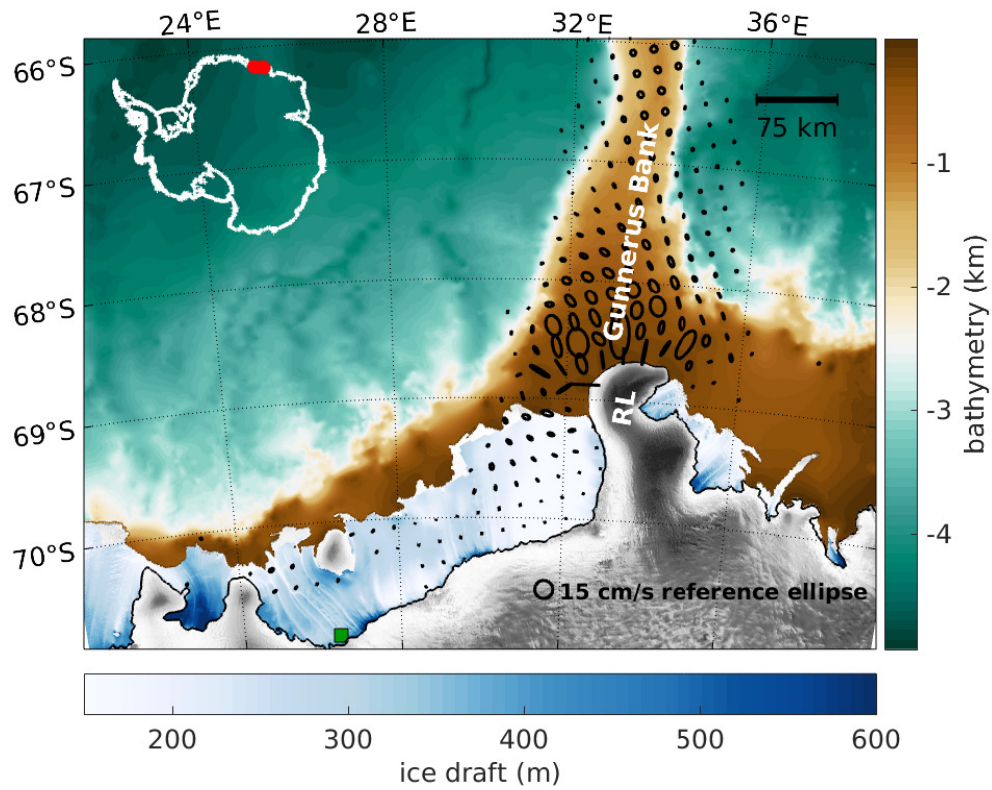
58 Sub-shelf melting depends on how fast energy can be transported across the bound-
59 ary layer to the ice-ocean interface. Jacobs *et al.* [1992] present three possible sources
60 that cause melting below Antarctic ice shelves. The first depends on the depression of
61 the local freezing point with increasing pressure and depth: even water that has been
62 cooled to the surface freezing temperature may supply heat to melt the base of an ice
63 shelf several hundred meters below sea level, driving the so called ice-pump circulation
64 [Lewis and Perkin, 1986]. The second source is warm deep water from the Southern Ocean,
65 which circulates along the continental shelf break. The third is the seasonally-warmed
66 Antarctic Surface Water (ASW), which could access the parts of the ice base close to the
67 calving front.

68 The delivery of energy for each of the three sources is modulated by processes in
69 the open ocean, as well as water circulation and transport in the ice shelf cavity [Din-
70 niman *et al.*, 2016; Stewart *et al.*, 2018]. The intrusion of warmer water into the ice shelf
71 cavities along the Dronning Maud Land coast is strongly controlled by the depth of the
72 Antarctic Slope Front thermocline [Hattermann, 2018] that is depressed below the shelf
73 break by wind forcing [Heywood *et al.*, 1998]. The on-shore heat transport across the front
74 is influenced by eddy overturning, freezing and melting of sea ice [Hattermann *et al.*, 2014;
75 Nøst *et al.*, 2011], and interactions with local topography [Dong *et al.*]. Persistent east-
76 erly winds along the entire coast of Dronning Maud Land (roughly between 30°E and
77 20°W) may push ASW into the ice shelf cavities [Zhou *et al.*, 2014], affecting the coastal
78 dynamics [Hattermann *et al.*, 2012], and potentially increasing sub-shelf melting [Hat-
79 termann *et al.*, 2014]. Mixing of Warm Deep Water and ASW produces Modified Warm

80 Deep Water (MWDW). Both ASW and MWDW are observed below the Fimbul Ice Shelf
81 [*Nøst et al.*, 2011; *Hattermann et al.*, 2012] as well as below the RBIS [*Callens*, 2014].

82 Ambient ocean waters inside the ice shelf cavity mix with melt water at the ice/ocean
83 boundary to form a buoyant plume that ascends in a boundary layer along the sloping
84 ice base [*Jenkins and Doake.*, 1991]. Tidal currents contribute to the mixing of the wa-
85 ter column and modify the hydrographic characteristics of the water masses at this bound-
86 ary [*Padman et al.*, 2018]. They also strongly strengthen the turbulence at the ice-ocean
87 interface, thereby increasing melting (or refreezing) [*Jourdain et al.*, 2019]. For the Filchner-
88 Ronne Ice Shelf, *Makinson et al.* [2011] show that, including tidal forcing in a numer-
89 ical ocean model leads to a three-fold increase in the modelled melt rate. Furthermore,
90 tides could trigger topographic waves over strong topographic slopes. They can be trapped
91 (e.g., if the generation site is poleward of the critical latitude, which is the case for di-
92 urnal tidal constituents in Antarctica) or resonate with the natural modes of basins or
93 bays depending on ocean and topography characteristics and excite vigorous internal waves
94 [*Semper and Darelius*, 2017; *Jensen et al.*, 2013].

95 The purpose of this study is to investigate the ice-ocean interaction mechanisms
96 driving sub-shelf melting for the RBIS, using an autonomous phase-sensitive radar sys-
97 tem (ApRES). The main advantage of this technique is the high vertical and temporal
98 resolution of vertical strain variations that can be obtained [*Nicholls et al.*, 2015], which
99 enables melt rates at a high temporal resolution. An additional advantage is that the
100 technique does not assume that the ice shelf is in steady state [*Corr et al.*, 2002]. Oceano-
101 graphic data from the ambient environment is collected to aid in interpreting the tem-
102 poral variability in melt rates. The observations give an insight into ice-ocean interac-
103 tion mechanisms that mediate sub-shelf melting along the Dronning Maud Land coast.
104 ApRES has been implemented in the other regions to derive time series of sub-shelf melt
105 rates [*Davis et al.*, 2018]. This is the first occasion that time series of sub-shelf melt rates
106 with such high temporal resolution have been obtained for this ice shelf. We explain the
107 observed variability in basal melt rates using directly measurements of temporal vari-
108 ability in the ocean.



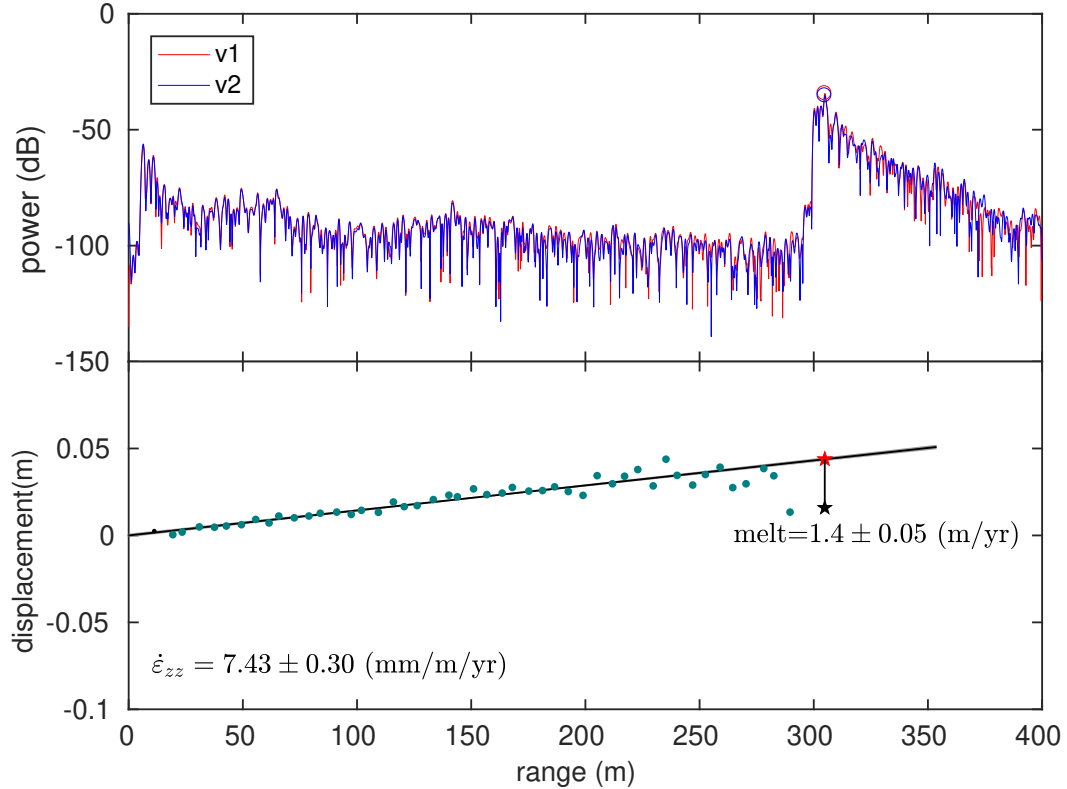
109 **Figure 1.** Overview of the study region. The background image is from a Radarsat mosaic
 110 [*Jezeq and RAMP-Product-Team, 2002*], overlaid with the ice draft [*Howat et al., 2019*]. The lo-
 111 cation of the ApRES on the RBIS is indicated by the green square. The grounding line is plotted
 112 in black [*Depoorter et al., 2013*]. Ocean bathymetry is from GEBCO_2014 [*Weatherall et al.,*
 113 2015], superimposed with tidal ellipses for the K1 constituent from the tide model CATS2008_opt
 114 where the amplitude of the semi-major axes is larger than 1.5 cm s^{-1} . RL denotes the Riiser-
 115 Larsen Peninsula.

116 2 Data collection and processing

117 2.1 Sub-shelf melt rates time series from radar measurements

118 The ApRES [Brennan *et al.*, 2014; Nicholls *et al.*, 2015] was deployed from Jan-
 119 uary to December in 2016 on the RBIS (Fig. 1) about 90 km from the ice-shelf front and
 120 5 km seaward from the grounded ice on the fast flowing portion of the West Ragnhild
 121 glacier, which is the third largest outlet glacier along the Dronning Maud Land Coast
 122 [Callens *et al.*, 2014]. The ice thickness at the site was ~ 300 m, but increases up to 600
 123 m in the grounding zone upstream, and ice flow velocities in this region range between
 124 250 and 300 m a⁻¹ [Rignot *et al.*, 2013]. By transmitting an electromagnetic signal and
 125 receiving the echo, the radar system can detect the ice base (ice-ocean interface) as well
 126 as relatively weak internal reflecting layers that are due to changes in ice permittivity
 127 (Fig. 2). Between two consecutive measurements, the relative vertical motion of inter-
 128 nal layers and the base can therefore be tracked. The displacements of the internal lay-
 129 ers determine how the thickness of the column evolves due to vertical strain and bottom
 130 melting (see below).

131 The equipment is described by Nicholls *et al.* [2015] and uses a Frequency Mod-
 132 ulated Continuous Wave technique. It retains the phase of the echo, ensuring a high pre-
 133 cision of the measurement, i.e., a signal-to-noise ratio of 17 dB allows a 1 degree change
 134 in phase to be detectable, corresponding to a 1 mm change in range in the case of a sys-
 135 tem centred on 300 MHz [Nicholls *et al.*, 2015]. The generated chirp frequency ranges
 136 from 200 MHz to 400 MHz. The bandwidth of 200 MHz gives a coarse range resolution
 137 of 43 cm in the vertical, and millimetre range precision can be achieved with relative phase
 138 measurements. A typical measurement lasts for 1-2 min. The instrument takes 20 mea-
 139 surements per hour to be averaged together, and sleeps between measurements. The strength
 140 of the signal decreases with depth, leading to a larger error in the phase tracking for the
 141 deeper reflectors. Therefore, reflectors deeper than 250 m are not used because of the
 142 noise. The result is a clean linear fit, which is the expected behaviour for a freely float-
 143 ing ice shelf. The gradient of the fitted line (Fig. 2) gives the average vertical strain rate
 144 from the non-compacting ice, which is nearly the whole column due to the lack of firm
 145 cover in this area [Lenaerts *et al.*, 2016]. With both spatial and temporal fluctuations
 146 in the vertical strain rate accurately quantified, it is possible to estimate the vertical dis-
 147 placement of the ice shelf base in response to ice flow divergence. Differences between



152 **Figure 2.** Example of a melt-rate calculation for ApRES based on two visits with a 7 day
 153 interval. The top panel shows radar records of the initial measurements (blue) and their cor-
 154 responding remeasurements (red). The circles label the reflection from the ice shelf base. The
 155 bottom panel shows the displacements of the reflecting surfaces and the base. We linearly fit
 156 (black line) the displacements (green dots) of reflecting surfaces and calculate the vertical strain
 157 rate. The black star marks the observed displacement of the shelf base while the red star marks
 158 the displacement due to vertical strain.

148 the predicted and observed motion of the basal reflector arise because of sub-shelf melt-
 149 ing or accretion (Fig. 2). All reflector displacements were processed using the method
 150 described by *Brennan et al.* [2014]. Details of the assumptions and derivations are given
 151 by *Jenkins et al.* [2006].

159 2.2 Sea-ice and ocean properties

160 Time series of sea-ice and ocean properties in this region are combined to deter-
 161 mine the mechanisms controlling the sub-shelf melting. Daily sea-ice concentration data
 162 in front of the ice shelf (24°– 34°E, 71°– 68°S) during 2016 are generated based on the

163 satellite passive microwave-derived data sets [Fetterer *et al.*, 2017]. Sea ice index data
 164 are provided by the National Snow and Ice Data Center.

165 Seasonal variations in the open ocean coastal hydrography are obtained from cli-
 166 matological data presented by Hattermann [2018], which is based on conductivity-temperature-
 167 depth (CTD) profiles from ships and Satellite Relay Data Logger-equipped seals (SRDL-
 168 CTD) [Boehme *et al.*, 2009] collected near the continental shelf break in the Kapp Norvegia
 169 region 10°– 25°E, 68°– 74°S in the period from 1977 to 2016 [Hattermann, 2018]. The
 170 full climatological dataset as well as the underlying raw CTD data are available via the
 171 pangaea repository [www.pangaea.de] and references therein Hattermann and Rohardt
 172 [2018]. Previous studies have shown the co-evolution of hydrographic properties along
 173 the Dronning Maud Land coast (e.g., [Nøst *et al.*, 2011]), and the few SRDL-CTD pro-
 174 files that exist north of the RBIS (Figs. S2 and S3) confirm that the Kapp Norvegia cli-
 175 matology is representative for the open ocean seasonality, despite being located further
 176 east. As an external forcing, temporal and spatial strength of tides in the study region
 177 are obtained using the CATS2008_opt barotropic tide model [Padman *et al.*, 2008].

178 Temporal variability of ocean properties below the RBIS are unknown. We there-
 179 fore present contemporaneously collected data from oceanographic moorings beneath the
 180 Fimbul ice shelf (M1 and M3 Hattermann *et al.* [2012]), which is located approximately
 181 1000 km further west along the Dronning Maud Land coast ($\sim 0^\circ\text{E}$) with a similar con-
 182 figuration and similar magnitude of melt rates [Rignot *et al.*, 2013] as the RBIS. The moored
 183 instruments are located at about 200 m depth, a few meters beneath the ice base.

184 **3 Temporal variability in sub-shelf melting and ocean properties**

185 Time series of basal melt rates (Fig. 3) are obtained for sliding intervals of both
 186 four hours and one day. That is, we compare pairs of observations separated by four hours
 187 and separated by one day, and then move that interval along by one hour.

188 The observational year 2016 can be divided into three periods corresponding to the
 189 phases of the melt rates (Fig. 3a): (a) From January to mid-May, the ApRES dataset
 190 indicates higher average melt rates of up to 10 m a^{-1} , as well as a higher variability in
 191 melt rates. This period coincides with the sea ice-free season (Fig. 3b). (b) From May
 192 to August, it changes abruptly to a phase of weak melt/refreezing when the sea ice cover
 193 increases at the onset of the winter. (c) From September, the magnitude and the vari-

194 ation of melt rates increase gradually. However, both are much lower than period (a),
 195 consistent with a significant sea ice cover during that period.

196 The evolution of open ocean stratification along the Dronning Maud Land coast
 197 is in accord with the three periods above (Fig. 3d, see also Fig. S2 for time series of tem-
 198 perature and salinity). Strong vertical density gradients in the upper 100-200 m are ob-
 199 served during the sea-ice-free period (a) when fresher and solar heated ASW accumu-
 200 lates along the coast [Zhou *et al.*, 2014]. At the onset of the sea ice formation period (b),
 201 the stratification abruptly vanishes when brine rejection convectively mixes the water
 202 column on the continental shelf. Later during winter (c) when further ice formation is
 203 suppressed by a solid ice cover that isolates the ocean from cold atmospheric temper-
 204 atures, along-shelf advection [Graham *et al.*, 2013] and on-shelf eddy fluxes of MWDW
 205 [Nøst *et al.*, 2011] slowly restratify the water column.

206 The Fimbul Ice Shelf moorings also show a seasonal inflow of solar heated surface
 207 water (Fig. 3e), albeit with a different timing because of the delay associated downwelling
 208 of the ASW before it enters the cavity [Zhou *et al.*, 2014]. However, the maximum tem-
 209 peratures observed around March and April and the gradual cooling until August/September
 210 does not fit well with the abrupt drop in melt rates seen at the ApRES site in mid-April.

211 To capture signals at higher frequencies (e.g. semi-diurnal), we use a sliding, 4-hour
 212 interval to obtain the melt rates for the wavelet analysis (Fig. 3a, c). A generalized Morse
 213 wavelet characterized by parameters $\gamma = 3$ and $\beta = 20$ is employed here [Lilly and
 214 Olhede, 2009]. Melt rate variability at the diurnal periods is dominant, and much larger
 215 than the semi-diurnal periods. However, vertical strain rates express a strong variabil-
 216 ity at both diurnal and semi-diurnal periods (not shown). During summer, there is a sec-
 217 ondary maximum in the 8-16 days range in melt rate variability, probably correspond-
 218 ing to the fortnight spring-neap cycle, which can also be seen from the pulsatile signal of
 219 the diurnal frequencies. For the whole observed period, the diurnal signal from the ApRES
 220 is strong in period (a), vanishes in period (b) and reappears weakly during period (c).

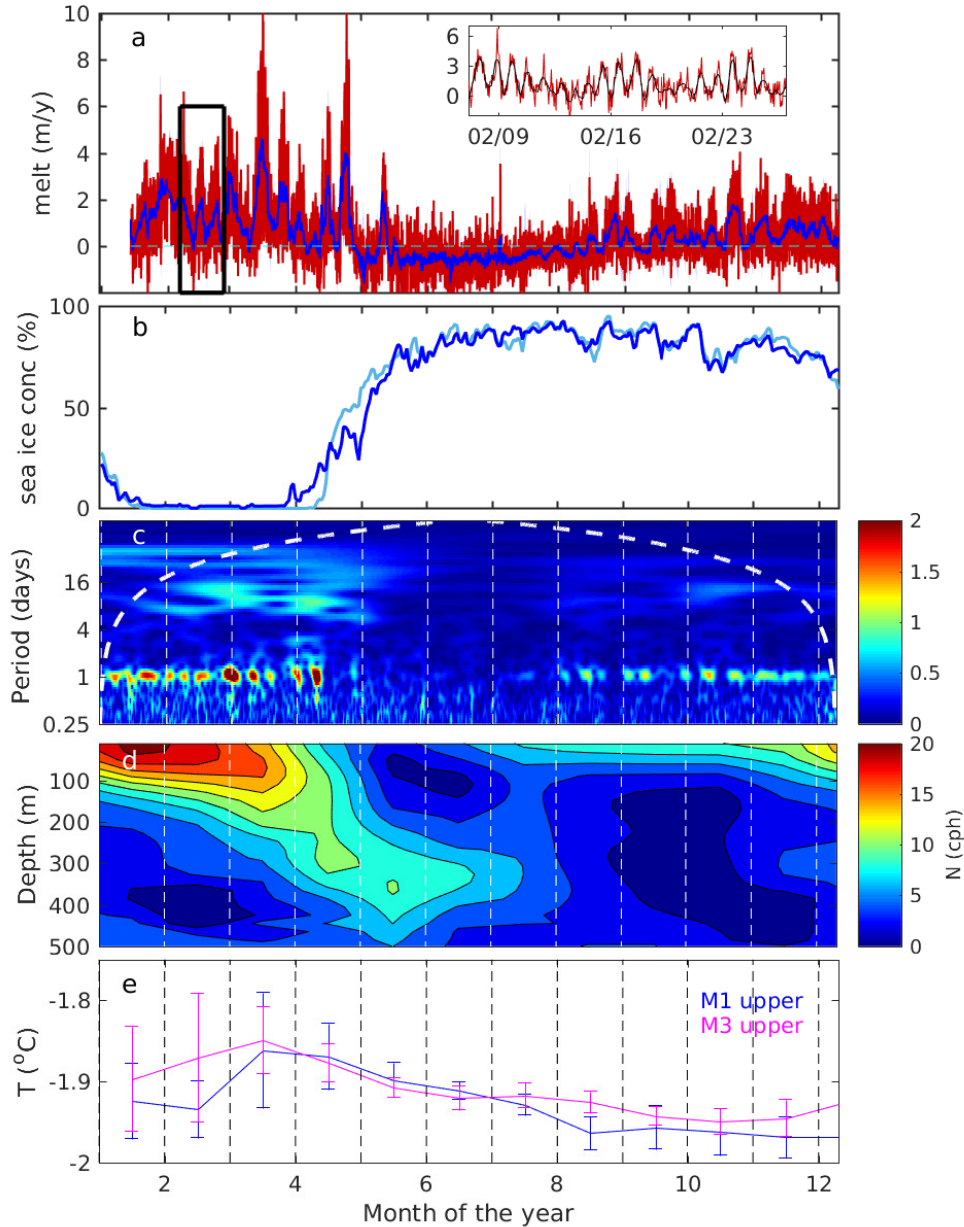
221 Figure. 1 also shows the dominant tidal forcing in the study region. Results from
 222 the CATS2008_opt model show a pronounced local amplification of diurnal tides over
 223 the Gunnerus Bank, with K1 tidal velocities of over 20 cm s^{-1} being an order of mag-
 224 nitude larger than generally found along the shelf break in the area (Fig. 1). Moreover,
 225 the second major diurnal constituent O1 shows a similar pattern, while amplitudes of

226 the semi-diurnal frequencies in contrast are much lower, with current velocities not ex-
227 ceeding a few cm per second and no pronounced amplification over the Gunnerus Bank
228 being seen in tidal model (Fig. S1). This asymmetry is attributed to the generation of
229 topographically trapped diurnal vorticity waves that have been reported in other regions
230 along other the Antarctic continental shelf break (e.g. [Middleton *et al.*, 1987; Padman
231 and Kottmeier, 2000; Padman *et al.*, 2003]), but also associated with divergent topog-
232 raphy as shown by [Skarohamar *et al.*, 2015].

244 4 Discussion

245 The coincidence of higher melt rates and the absence of sea ice indicates a nearly
246 instantaneous link between the melt rates and the seasonal changes in the oceanic en-
247 vironment seaward of the ice front. Downwelling of solar-heated ASW in the sea ice-free
248 period has been observed to enhance basal melting [Hattermann *et al.*, 2012, 2014] and
249 models suggest that this mechanism is a general feature of the narrow continental shelf
250 configuration with easterly winds [Zhou *et al.*, 2014]. However, it is questionable that
251 this process will affect the deeper parts of the cavity to instantaneously increase melt
252 rates near the grounding line. ASW is usually formed in a relatively thin surface layer
253 when sea ice melts. It gradually spreads beneath ice shelf with much of the warming sig-
254 nal being lost on its way into the cavity, as also be seen from the rather moderate tem-
255 perature increase at the Fimbul moorings (Fig. 3e), which are both shallower and closer
256 to the open ocean than the ApRES measurements in this study obtained at 300 m depth
257 90 km away from the ice shelf front.

258 Instead of relating the melt rates to the lateral advection of warmer water, we sug-
259 gest that the observed variability is controlled by the seasonal presence of baroclinic waves
260 inside the cavity that modulate the vertical heat flux toward the ice base. In this mech-
261 anism, topographically trapped vorticity waves generated by diurnal tides over sloping
262 topography provide energy to excite internal waves that propagate along the slopes of
263 the grounding line toward the ApRES site. The melt rate variability observed beneath
264 RBIS closely follows the upper open ocean stratification and the wavelet analysis indi-
265 cates that the local circulation is more energetic at diurnal than at semi-diurnal frequen-
266 cies (Fig. 3c). Together with the strong diurnal tidal forcing over the Gunnerus Bank
267 (Fig. 1), this supports the hypothesis of seasonal resonance of internal waves beneath
268 the ice shelf. In particular the abrupt termination of the diurnal melt rate signal at the



233 **Figure 3.** Time series of the observational datasets: (a) time series of melt rates with 4-hour
 234 window (red line), overlaid by time series with one-day window (blue line). The time series of
 235 melt rates with 4-hour window from 7th February to 26th February (the black box) is shown in
 236 the inset plot (red line), overlaid by the smoothed data (black line) to clearly show the diurnal
 237 period; (b) sea ice concentration at the front of the RBIS in 2016 between 24° and 34°E in cyan
 238 and between 30° and 32°E in blue; (c) periodograms of the melt rates derived using 4-hour mov-
 239 ing window. The white dashed line shows the cone of influence; (d) climatological time series of
 240 coastal ocean stratification seaward of the ice front at Kapp Norvegia [Hattermann, 2018]; (e)
 241 time series (monthly mean) of 2016 upper ocean temperatures at two moorings beneath the Fim-
 242 bul Ice Shelf (near the ice front) at the same location as in Hattermann *et al.* [2012]. The error
 243 bars represent the variance within each month of the hourly data (standard deviation).

269 onset of sea ice formation stands out and is hard to explain by other mechanisms than
270 a ceasing of those waves as convection mixes the water column, before a gradual restrat-
271 ification during winter allows baroclinic waves to be generated and transmitted once again.

272 Unlike freely propagating, semi-diurnal modes, the diurnal topographic waves as-
273 sociated with the barotropic tides over the Gunnerus Bank are trapped, i.e. they can-
274 not propagate, because they are poleward of the critical latitude, and their conversion
275 into internal waves may be important to facilitate the resonance of these modes inside
276 the ice shelf cavity. Non-linear and linear examples of propagating baroclinic waves at
277 above-critical latitudes exist [*Rippeth et al.*, 2017; *Hughes and Klymak*, 2019] and com-
278 parable seasonal resonances of baroclinic waves with diurnal tides have been observed
279 at similar configurations along the Antarctic continental slope [*Semper and Darelius*, 2017;
280 *Jensen et al.*, 2013]. Idealized models that have been used to quantify the internal wave
281 energy rely on constant geometry assumptions, which is challenging for the complex and
282 partially unknown bathymetry around the Riiser-Larsen Peninsula peninsula and beneath
283 the RBIS. However, *Falahat and Nycander* [2015] find that Gunnerus Bank experiences
284 some of the highest energy densities of bottom-trapped internal tides around the Antarc-
285 tic continent. Furthermore, the conversion into internal waves may be important for sur-
286 passing the strong potential vorticity barrier that is imposed to a barotropic ocean by
287 the 100 m to 200 m step change in watercolumn thickness at the quasi-vertical calving
288 front of the ice shelf that hampers the exchange of the cavity circulation with the open
289 ocean [*Nicholls et al.*, 2009]. Vertical density gradients at the depth of the ice shelf base
290 may decouple the lower part of the water column, effectively weakening the potential vor-
291 ticity barrier and allowing baroclinic signals to propagate from the open ocean into the
292 cavity. Such decoupling by summer time stratification has been evoked to explain the
293 exchange across ice fronts at other ice shelves [*Nicholls et al.*, 2009; *Darelius and Sallée*,
294 2018; *Malyarenko et al.*, 2019]. With wavelengths usually in the order of a few hundred
295 kilometers for the first modes, the resulting internal waves may well resonate with the
296 size of the ice shelf cavity, causing a nearly instantaneous response to open ocean sig-
297 nals along the grounding line from the outcrop to the open ocean at the Riiser-Larsen
298 Peninsula toward the ApRES.

299 The observed basal melt rates are essentially controlled by the local thermal driv-
300 ing and the friction velocity [*Holland and Jenkins*, 1999]. The presence of energetic waves
301 would increase both of these two factors by enhancing the upward mixing of heat in the

302 ambient ocean and by increasing turbulence at the ice-ocean interface. Herein, the ob-
303 served magnitude of the summer peak melt rates suggests that a significant amount of
304 water somewhat warmer than the surface freezing point is available in the cavity to be
305 raised by tidal activity. While the bathymetry beneath RBIS is largely unknown, warmer
306 inflows at depth may also affect the melt rate variability. *Hattermann* [2018] suggested
307 that the presence of ASW along the coast of Dronning Maud Land causes a shoaling of
308 the thermocline along the shelf break that may cause seasonal access of warmer water
309 over sills and troughs into the ice shelf cavity. Although such intermittent inflows prob-
310 ably play a role in providing ocean heat for melting inside the cavity, large variations in
311 this process would likely leave more gradual and somewhat delayed signature in the melt
312 rate signal at the grounding line than the observed variability that appears to be phase
313 locked with the sea-ice-free period as is consistent with the topographic wave argument.

314 Finally, if tides and topographic waves beneath RBIS would have the same char-
315 acteristics all year round, the wavelet analysis in Fig. 3c would show an enhancement
316 of the diurnal frequency interval also during the sea ice formation period, as melt rate
317 variability in Fig. 3a is not zero at that time. Instead, most of the variability seems to
318 occur at higher frequencies, suggesting the absence of a pronounced diurnal signal in the
319 exchange velocities at the ice-ocean boundary.

320 5 Conclusions

321 Based on continuous ApRES measurements, we provide a yearlong, hourly time se-
322 ries of directly measured sub-shelf melt rates that enable us to investigate temporal vari-
323 ability of sub-shelf melt rates at a broad range of time scales.

324 The magnitude of sub-shelf melt rates varies from close to 0 in winter to 10 m a^{-1}
325 in summer, with variations over a broad range of time scales, from tidal to seasonal. We
326 propose that the sub-shelf melt rate of the RBIS near its grounding line is controlled by
327 topographic waves triggered over the Gunnerus Bank. By controlling the turbulent mix-
328 ing of heat and salt towards the ice shelf base, topographic waves directly affect the lo-
329 cal melt rates and act as a conduit for propagating open ocean changes to the ground-
330 ing line far inside the ice shelf cavity, even without advecting different water masses be-
331 neath the ice shelf. Such topographic waves are observed to resonate with the diurnal
332 tides at the Antarctic continental slope [*Semper and Darelius, 2017*], but their resonance

333 depends on the stratification in the upper water column, leading to a seasonal variation
334 in the strength of waves, which, in turn, leads to the seasonally varying strength of the
335 diurnal signal in the melt rate time series. The formation and decline of sea ice there-
336 fore impacts the melt rates indirectly by modulating the stratification of the ocean.

337 Although the detailed dynamics of the internal waves inside the RBIS cavity will
338 need to be explored in future studies, our observations suggest a mechanism for regu-
339 lating ice shelf basal melting that has not previously been considered. It seems to dom-
340 inate the melt variability along parts of the grounding line at the RBIS and will likely
341 be present beneath other ice shelves along the Dronning Maud Land coast. In this pro-
342 cess, open-ocean bathymetry and far-field tidal and surface forcing plays a key role in
343 controlling melt variability deep inside the cavity. It may also be important for modu-
344 lating the melt rate response to future climate change, as observations and models sug-
345 gest an ongoing freshening of ASW [*de Lavergne et al.*, 2014]. Recent studies based on
346 fully coupled ice-ocean-atmospheric modelling show that the melt water flux from Antarc-
347 tica will trap warm water below the sea surface, in turn increase the melting near the
348 grounding line and create a positive feedback [*Golledge et al.*, 2019; *Menviel et al.*, 2010;
349 *Bronselaeer et al.*, 2018]. Increasing melt-water fluxes will further increase the (winter-
350 time) stratification of continental shelf waters, and in case of the RBIS, this will likely
351 enhance the presence of topographic waves and permanently increase the heat flux and
352 melt rates near the grounding line of the ice shelf.

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362 bul Ice Shelf mooring data, which are available via <https://data.npolar.no> upon re-
363 quest.

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 367 W. Rack, G. Rotschky, S. Urbini, P. Vornberger, and N. Young (2011), Getting
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