

Tanguy Mertens

A new mapped infinite partition of unity  
method for convected acoustical radiation in  
infinite domains.

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## Remerciements

*Si tu donnes un poisson à un homme, il ne mangera qu'un jour. S'il apprend à pêcher, il mangera toute sa vie.*

Proverbe de Confucius, repris plus tard par Dominique Pire dans le cadre de l'action Iles de Paix.

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Enfin, je terminerai par remercier les deux petits rayons de soleil qui ont ensoleillé la fin de ce long parcours.

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## List of Symbols

### Greek symbols

|                    |                                                                   |               |
|--------------------|-------------------------------------------------------------------|---------------|
| $\beta$            | : $\sqrt{1 - M_0^2}$                                              |               |
| $\gamma$           | : Poisson ratio of specific heat capacities : $c_p/c_v$           |               |
| $\Gamma$           | : interface separating the inner and the outer domains            |               |
| $\varepsilon$      | : Error                                                           |               |
| $\mu$              | : phase function                                                  | $[m]$         |
| $\rho$             | : mass density                                                    | $[kgm^{-3}]$  |
| $\rho_0$           | : steady mean density                                             | $[kgm^{-3}]$  |
| $\rho_a$           | : acoustic density                                                | $[kgm^{-3}]$  |
| $\sigma$           | : stress tensor                                                   | $[Nm^{-2}]$   |
| $\phi$             | : velocity potential                                              | $[m^2s^{-1}]$ |
| $\phi_0$           | : mean velocity potential                                         | $[m^2s^{-1}]$ |
| $\phi_a$           | : acoustic velocity potential                                     | $[m^2s^{-1}]$ |
| $\tilde{\phi}_a$   | : amplitude of the harmonic acoustic velocity potential           | $[m^2s^{-1}]$ |
| $\tilde{\phi}^h$   | : numerical approximation of $\tilde{\phi}_a$                     | $[m^2s^{-1}]$ |
| $\tilde{\phi}_h^I$ | : numerical approximation in the outer region $\Omega_o$          | $[m^2s^{-1}]$ |
| $\Phi_\alpha$      | : shape function for the $\alpha^{th}$ degree of freedom          |               |
| $\Phi_\alpha^I$    | : infinite shape function for the $\alpha^{th}$ degree of freedom |               |
| $\omega$           | : angular frequency                                               | $[s^{-1}]$    |
| $\Omega$           | : domain                                                          |               |
| $\Omega_i$         | : inner region                                                    |               |
| $\Omega_o$         | : outer region                                                    |               |

## Arabic symbols

|                  |                                                     |                    |
|------------------|-----------------------------------------------------|--------------------|
| $\tilde{a}_n$    | : normal acceleration of a vibrating wall           | $[ms^{-2}]$        |
| $A_n$            | : normal acoustic admittance                        | $[m^2 skg^{-1}]$   |
| $A_{mn}^{\pm}$   | : incident and reflected modal amplitude            | $[m^2 s^{-1}]$     |
| $c$              | : speed of sound                                    | $[ms^{-1}]$        |
| $c_0$            | : steady mean part of the speed of sound            | $[ms^{-1}]$        |
| $c_{\infty}$     | : speed of sound at large distance from the source  | $[ms^{-1}]$        |
| $c_p$            | : specific heat capacity at constant pressure       | $[JK^{-1}]$        |
| $c_v$            | : specific heat capacity at constant volume         | $[JK^{-1}]$        |
| $dofs$           | : number of unknowns of the approximation           |                    |
| $E$              | : energy flow out of a surface                      | $[J]$              |
| $E_{mn}^{\pm}$   | : incident and reflected modal pattern              |                    |
| $f$              | : excitation frequency                              | $[s^{-1}]$         |
| $G$              | : geometric factor                                  |                    |
| $h$              | : mesh size                                         | $[m]$              |
| $H$              | : Hilbert space                                     |                    |
| $i$              | : imaginary unit = $\sqrt{-1}$                      |                    |
| $\mathbf{I}$     | : Sound intensity                                   | $[Wm^{-2}]$        |
| $J'$             | : stagnation entropy                                | $[Jkg^{-1}]$       |
| $k$              | : wavenumber                                        | $[m^{-1}]$         |
| $k_{r,mn}^{\pm}$ | : incident and reflected radial wavenumber          | $[m^{-1}]$         |
| $k_B$            | : Boltzmann constant                                | $[JK^{-1}]$        |
| $K_{z,mn}^{\pm}$ | : incident and reflected axial wavenumber           | $[m^{-1}]$         |
| $L_j^d$          | : Legendre polynomial of order $d$ for node $j$     |                    |
| $L_s$            | : curve enclosing the boundary $S_s$                |                    |
| $L_v$            | : curve enclosing the boundary $S_v$                |                    |
| $m$              | : angular mode number                               |                    |
| $\mathbf{m}'$    | : mass flux                                         | $[kgm^{-2}s^{-1}]$ |
| $m_0$            | : radial order of the infinite element              |                    |
| $m_w$            | : mass of a molecule                                | $[kg]$             |
| $M_0$            | : mach number                                       |                    |
| $M_i$            | : Mapping function for node/point $i$               |                    |
| $\mathbf{n}$     | : outer normal to the domain                        |                    |
| $n$              | : radial mode number                                |                    |
| $n_d^I$          | : number of infinite degree of freedom              |                    |
| $n(j)$           | : size of the local approximation space at node $j$ |                    |
| $nni$            | : number of infinite nodes                          |                    |
| $nodes$          | : number of nodes                                   |                    |
| $N_i$            | : Partition of Unity function of node $i$           |                    |
| $N_m$            | : number of angular modes                           |                    |
| $N_n$            | : number of radial modes                            |                    |
| $N_M$            | : number of reflected modes (unknown)               |                    |

|                        |                                                                   |                     |
|------------------------|-------------------------------------------------------------------|---------------------|
| $p$                    | : fluid pressure                                                  | $[Pa]$              |
| $p_0$                  | : steady mean fluid pressure                                      | $[Pa]$              |
| $p_a$                  | : acoustic pressure                                               | $[Pa]$              |
| $\tilde{p}_a$          | : amplitude of the harmonic acoustic pressure                     | $[Pa]$              |
| $\tilde{p}_{an}$       | : analytic amplitude of the harmonic acoustic pressure            | $[Pa]$              |
| $\mathbf{q}$           | : heat flux                                                       | $[Wm^{-2}]$         |
| $Q_w$                  | : heat production                                                 | $[J]$               |
| $r_o$                  | : distance to the source point                                    | $[m]$               |
| $R$                    | : specific gas constant                                           | $[JK^{-1}mol^{-1}]$ |
| $R_j$                  | : radial function for infinite node $j$                           |                     |
| $R_j^d$                | : radial function of order $d$ for node $j$                       |                     |
| $s$                    | : entropy                                                         | $[Jkg^{-1}K^{-1}]$  |
| $S$                    | : boundary                                                        |                     |
| $S_i$                  | : mapping functions for the interface $\Gamma$                    |                     |
| $S_M$                  | : Modal boundary                                                  |                     |
| $S_s$                  | : soft wall                                                       |                     |
| $S_v$                  | : vibrating wall                                                  |                     |
| $t$                    | : time                                                            | $[s]$               |
| $T$                    | : Temperature                                                     | $[K]$               |
| $T_j$                  | : circumferential function for infinite node $j$                  |                     |
| $\tilde{u}_n$          | : normal displacement of a vibrating wall                         | $[m]$               |
| $\mathbf{v}$           | : fluid velocity                                                  | $[ms^{-1}]$         |
| $\mathbf{v}_0$         | : steady mean fluid velocity                                      | $[ms^{-1}]$         |
| $\mathbf{v}_\infty$    | : fluid velocity at large distance from the source                | $[ms^{-1}]$         |
| $\mathbf{v}_a$         | : acoustic velocity                                               | $[ms^{-1}]$         |
| $\tilde{\mathbf{v}}_a$ | : amplitude of the harmonic acoustic velocity                     | $[ms^{-1}]$         |
| $\mathcal{V}$          | : the Sobolev space $W^{1,2} = H^1 = \{f : f, \nabla f \in L^2\}$ |                     |
| $V_{jl}$               | : $l^{th}$ local approximation function of node $j$               |                     |
| $\tilde{w}_n$          | : normal velocity of a vibrating wall                             | $[ms^{-1}]$         |
| $W_j$                  | : weight function of node $j$                                     |                     |
| $W_j^I$                | : infinite weight function of the infinite node $j$               |                     |
| $W_{M,nm}$             | : modal weight function of the angular and radial mode $(m, n)$   |                     |

## Operators

|                   |                                         |
|-------------------|-----------------------------------------|
| $\nabla$          | : gradient operator                     |
| $\nabla \cdot$    | : divergence operator                   |
| $\nabla \times$   | : curl operator                         |
| $\Delta$          | : Laplacian operator                    |
| $\frac{D}{Dt}$    | : Total time derivative                 |
| $:$               | : the double dot product of two tensors |
| $\langle \rangle$ | : time average                          |
| $\Re$             | : Real part                             |



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## Industrial application: Turbofan radiation

The previous chapters verify the Mapped Infinite Partition of Unity formulation and illustrate the performances of the method by analysing several academic applications. This chapter considers the radiation of a turbofan engine.

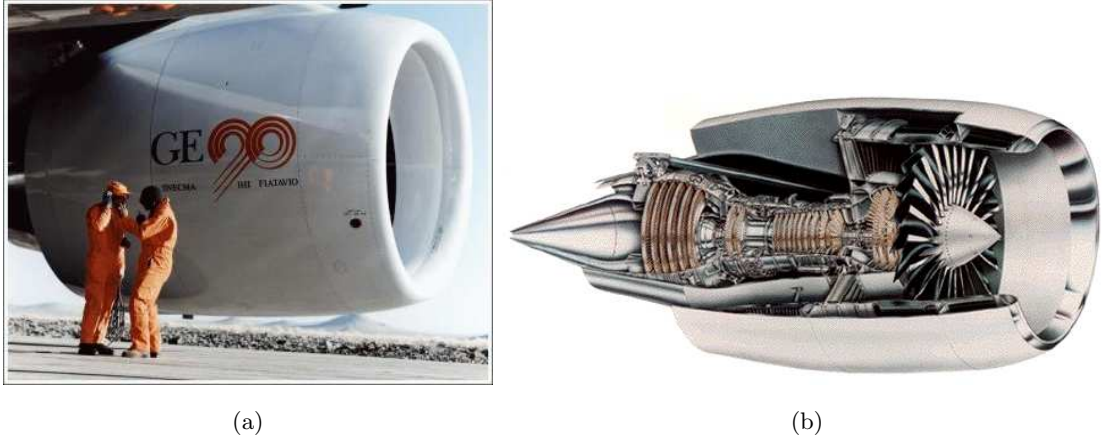
Aircraft noise is a major concern for approximately 100 square kilometers surrounding most major airports [81]. Stringent registrations tends to limit this environmental noise in spite of traffic growth which is planned to double in less than 15 years (India and China air traffic is expanding at 40% per year). There is then huge pressure on airplane manufacturers to significantly reduce aircraft noise.

A lot of research has already been conducted to reduce aircraft noise. The european project SILENCE(R) [80], for instance, was launched in 2001 with the aim to reduce the noise by up to 6 decibels as of 2008. It involved 51 companies and was funded with a total budget which exceeding 110 M Euro.

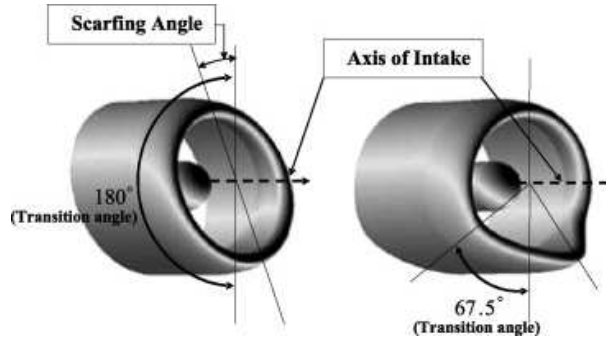
Aircraft noise is composed of three principal sources: airframe, jet and fan. The airframe is aero-acoustic noise corresponding to turbulent flow. This airframe noise is generated by landing gear and high lift devices. Engine noise is composed of jet noise which is emitted at the back of the engine and fan noise which propagates towards the inlet before being radiated in front of the engine (figure 8.1).

We decided to focus on the noise radiated by the fan because high bypass ratio has led to significant reduction in jet noise. The fan noise is composed of tonal and broadband noise. The reduction in fan noise can be obtained by acting on the source (design of low noise compressors) or on the propagation through the inlet. The perceived noise can be reduced by attenuation through acoustic treatment in the nacelle (liner) or by diverting the noise away from the ground with scarf or scoop inlets (figure 8.2). Our goal is to simulate fan propagation through the inlet and compute the radiated pressure and evaluate the effect of acoustic treatment or the influence of the inlet geometry.

The noise radiated by the fan may be separated in several contributions. The tonal noise corresponds to rotor alone or rotor-stator interactions and the broadband noise is due to turbulence-fan interactions. The rotor alone noise has a frequency corresponding to



**Fig. 8.1.** General Electric GE90-115B high bypass turbofan [83].



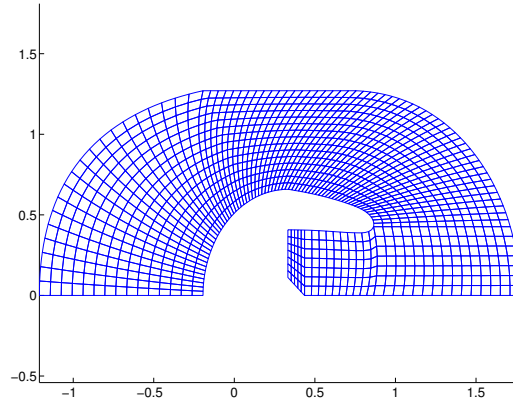
**Fig. 8.2.** Illustration of scarf (left) and scoop (right) inlet [82].

the blade passing frequency (BPF) and is represented by the first radial mode of the  $B^{th}$  azimuthal mode  $(m, n) = (B, 1)$ , where  $B$  is the number of blades. At subsonic fan speed, the rotor alone noise is cut-off. This means that it is an evanescent mode and it does not radiate noise. However, during takeoff, the fan speed at the tip of the blades is supersonic and the rotor alone noise is cut-on, it has then to be taken into account in simulations. In the case of rotor-stator interactions, rotor wakes impinge on stator vanes. The dominant tones correspond to harmonics of the blade passing frequency. According to the theory of Sofrin and Tyler [78], the corresponding azimuthal mode is given by the following relation:  $m = iB \pm jV$ , where  $V$  is the number of stator vanes,  $i$  is the time harmonic index and  $j$  is an integer number. Note that tonal noise is prescribed with a small number of modes while broadband noise is modelled by a multi-mode source of fixed energy.

A full nacelle analysis requires three-dimensional modelling. Unfortunately such applications are computationally demanding. A first approach consists in simulating nacelle radiation by an axisymmetric model which can take into account admittance and non-uniform flow but only considers axisymmetric geometries (no scarf, no spliced acoustic liners). Note that Duta and Giles [79] described a spectral formulation allowing to decompose the circumferential geometry by a Fourier spectral representation. They reduced a

three-dimensional simulation of an asymmetry inlet with a scarfing angle of about  $5^\circ$  with approximately 3 million of unknowns to a total of 0.24 million of unknowns for the spectral approximation.

The Mapped Infinite Partition of Unity Method has been developed for axisymmetric geometry as well as three-dimensional applications. We illustrate in this chapter nacelle radiation with the axisymmetric formulation. The geometry is realistic in the sense that it approximates the dimensions and curvatures of modern turbofan excepted for non-axisymmetric features (azimuthally segmented liners, scarf or scoop angle). The geometry of the application is illustrated with a mesh of 1260 elements in figure 8.3. We analyse the radiation of the rotor alone tone of an engine of 26 fan blades operating close to its maximum speed corresponding to the  $(26, 1)$  mode at blade passing frequency: nondimensional wave number  $ka = 30$  ( $f = 4000\text{Hz}$  and  $a = 0.4\text{m}$ ). These characteristics aim to represent real applications and are inspired from an analysis performed by Astley [84].



**Fig. 8.3.** Illustration of the geometry of the nacelle and an associated inner mesh of 1260 Q8 quadratic elements. Note that infinite elements (not represented) are attached to the interface  $\Gamma$  with the outer region.

The current analysis illustrates the radiation of mode  $(26, 1)$  at  $ka = 30$  for three different cases. We first consider no flow propagation, then take into account non-uniform flow and finally apply acoustic treatment within the convected simulation. The radiation of each case is illustrated with plots of the pressure within the computational domain. We also use directivity patterns to select appropriate discretizations and compare the constant enrichment, second order enrichment of the Mapped Infinite Partition of Unity Method and quadratic Finite Element simulations (ACTRAN<sup>TM</sup>). The infinite radial order used for the computations is  $m_0 = 15$ .

## 8.1 Radiation without flow

The distribution of the pressure is illustrated in figure 8.4 for the inner region and a partly in the outer region.

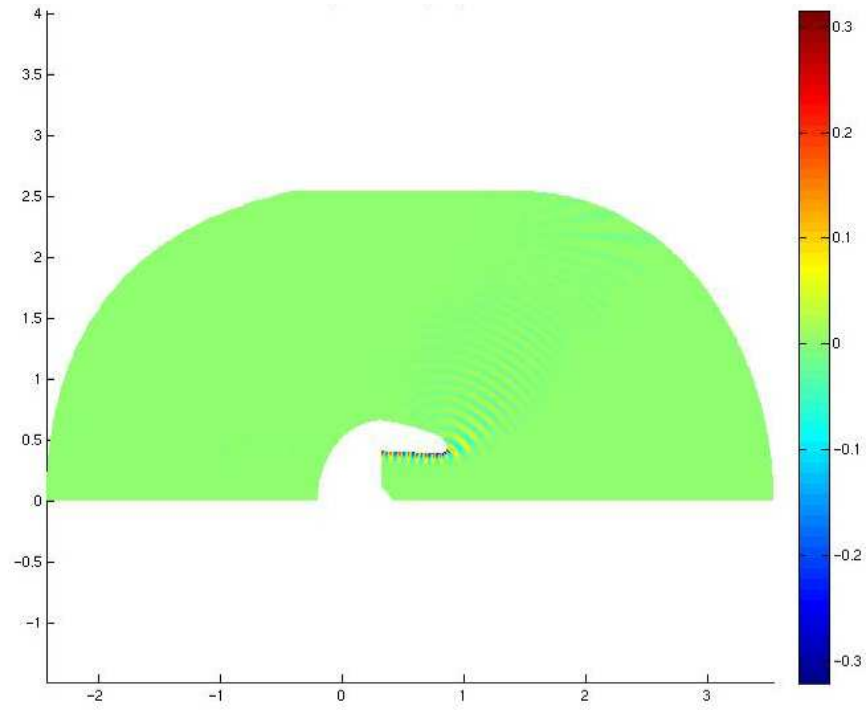
We compare directivity plots obtained with the Mapped Infinite Partition of Unity Method. The aim is to find the appropriate discretization for the application and compare the required number of degrees of freedom for the constant enrichment and the second order enrichment.

Directivity patterns (figure 8.5) represent the pressure in decibels<sup>1</sup> for a point located on a circle of radius  $r = 2m$  and centered at the origin. Note that the maximal pressure in the nacelle is over  $142dB$ . We observe that the Mapped Infinite Partition of Unity Method with second order enrichment converges already with 9000 degrees of freedom while the constant enrichment requires 36000 degrees of freedom.

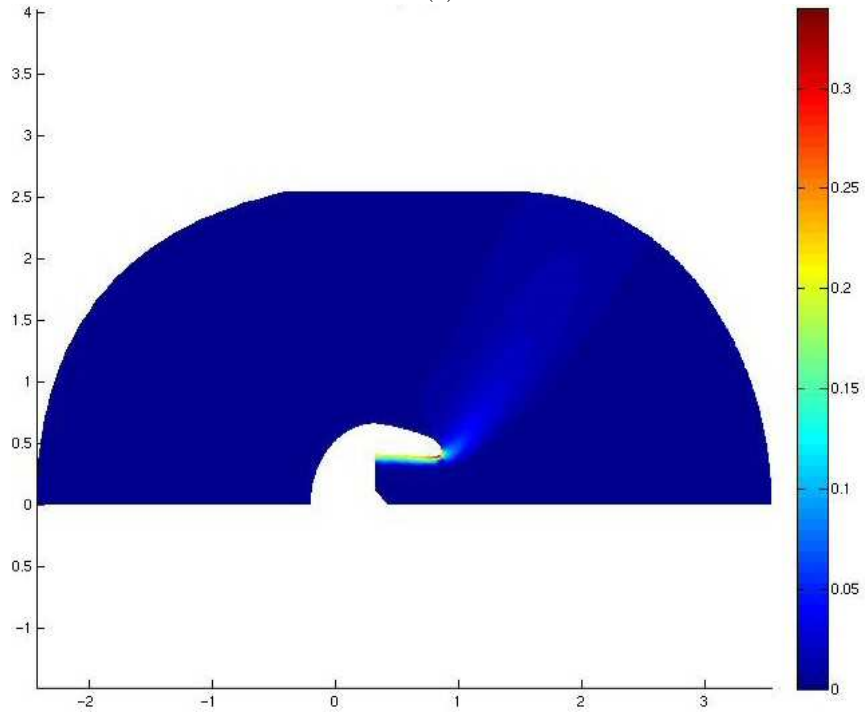
We also checked the validity of using an infinite radial order of  $m_0 = 15$ . Figure 8.6 illustrates the directivity pattern computed for three different radial orders:  $m_0 = \{10, 15, 20\}$ . The radial order  $m_0 = 10$  gives wrong results. The two other radial orders show close results, excepted discrepancies at angles close to  $80^\circ$ .

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<sup>1</sup>  $p_{dB} = \left(20 \log \left( \frac{p}{p_{ref}} \right) \right)$  with  $p_{ref} = 0.00002 Pa$

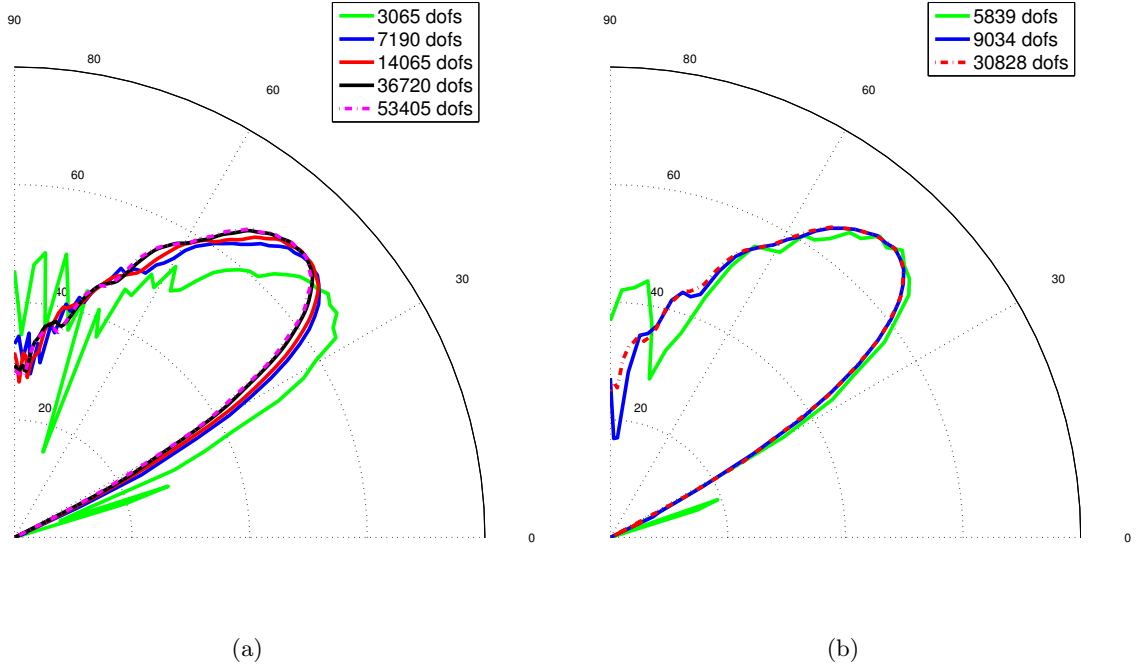


(a)

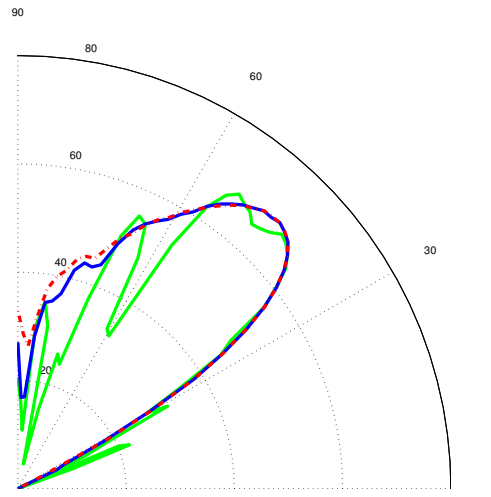


(b)

**Fig. 8.4.** MIPUM 0-15-1 (36720 degrees of freedom): Radiation of mode (26,1) at  $ka = 30$  without flow: (a) real part of the pressure (b) absolute value of pressure.



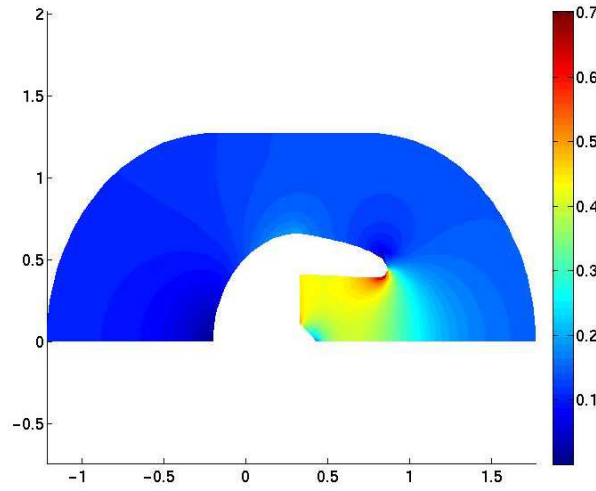
**Fig. 8.5.** Directivity pattern for the radiation of mode (26, 1) at  $ka = 30$  without flow: (a) MIPUM 0-15-1 (b) MIPUM 2-15-2.



**Fig. 8.6.** Directivity pattern for the radiation of mode (26, 1) at  $ka = 30$  without flow: MIPUM 2-10-2 (—), MIPUM 2-15-2 (—) and MIPUM 2-20-2 (—).

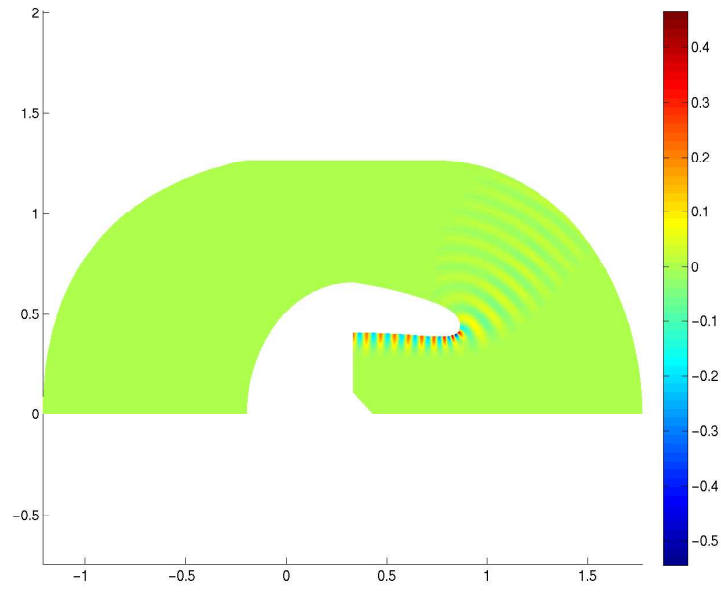
## 8.2 Convected radiation

The convected radiation is illustrated in figure 8.8. The input steady mean flow is illustrated in figure 8.7 for the inner region. In the outer region, the flow is assumed to be uniform and oriented in the axial direction:  $\mathbf{v}_0 = -50 \mathbf{1}_z m/s$ . The mean flow has been computed by the software *ACTRAN<sup>TM</sup>*. Note that the mean flow is computed on a mesh of 2080 nodes and then interpolated on the nodes of the acoustic mesh.

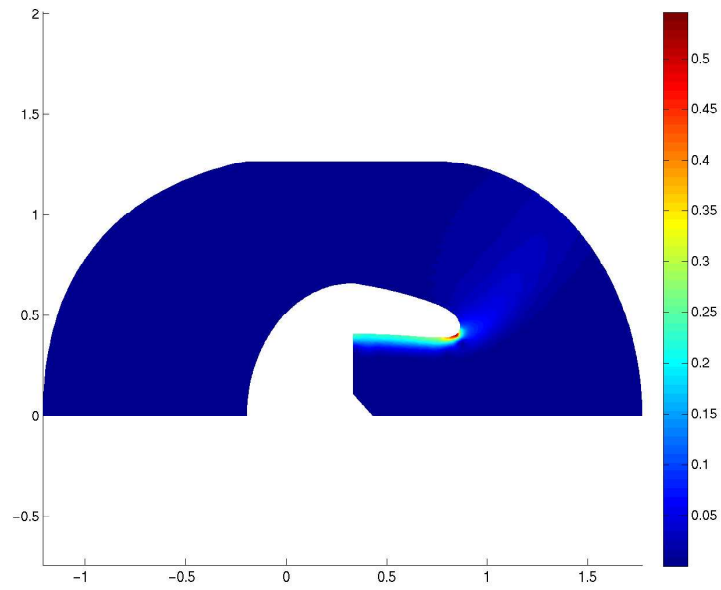


**Fig. 8.7.** Illustration of the steady mean flow distribution: Mach number

The directivity pattern at a distance of  $2m$  for the convected radiation is shown at figure 8.9 for different discretizations. The maximal value of the pressure in the nacelle is about  $138dB$ . Figure 8.10 compares directivity for converged solutions. This shows that both MIPUM 2-15-2 and quadratic Finite Element Method converges towards the same solution.

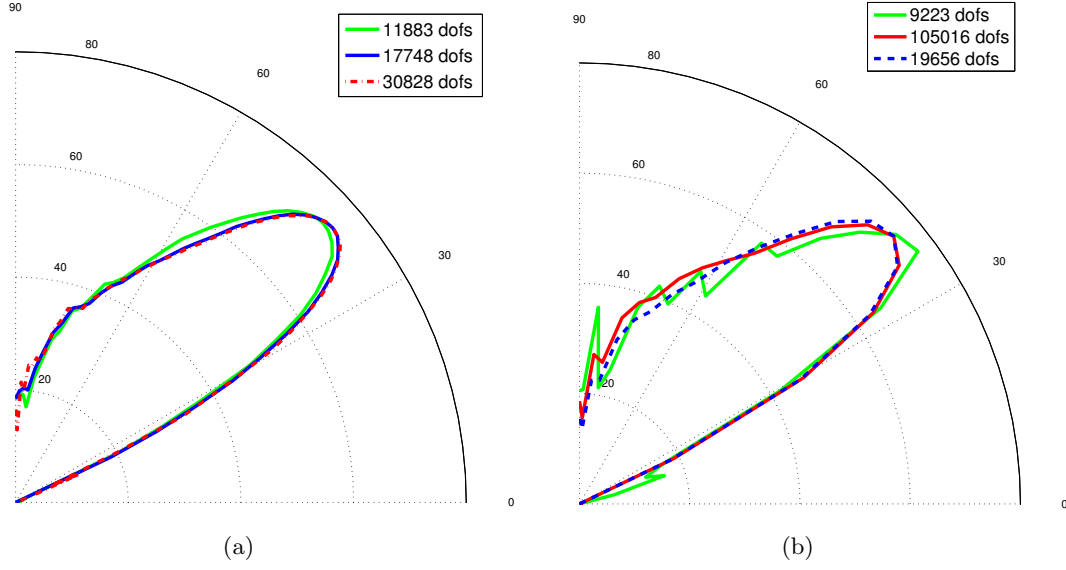


(a)

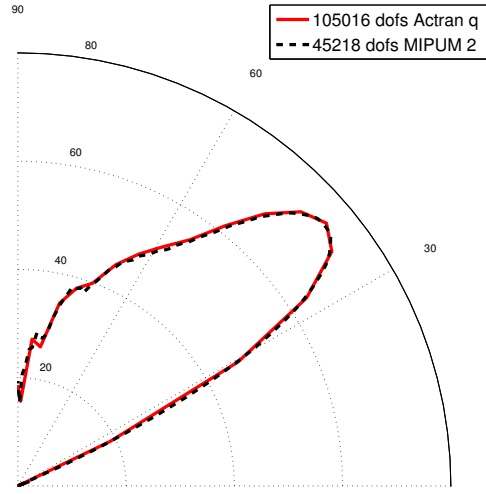


(b)

**Fig. 8.8.** Quadratic Finite Element Method (105016 degrees of freedom): Radiation of mode  $(26, 1)$  at  $ka = 30$  without acoustic treatment: (a) real part of the pressure (b) absolute value of pressure.



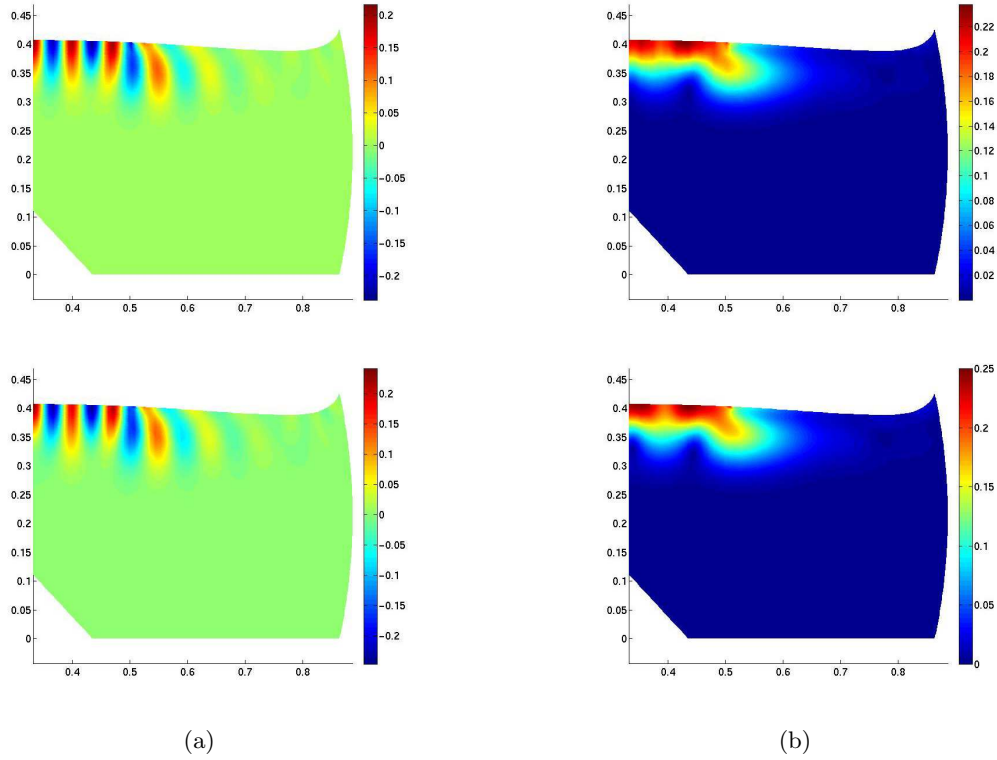
**Fig. 8.9.** Directivity pattern for the convected radiation of mode  $(26, 1)$  at  $ka = 30$  without acoustic treatment: (a) MIPUM 2-15-2 (b) QFEM.



**Fig. 8.10.** Directivity pattern for the convected radiation of mode  $(26, 1)$  at  $ka = 30$  without acoustic treatment: comparison of converged directivity for QFEM and MIPUM 2-15-2.

### 8.3 Convected radiation and influence of liners

A liner is inserted in the nacelle from  $x = 0.5m$  to  $x = 0.75m$ . It is simulated by prescribing an admittance boundary condition with an admittance of:  $A_n = 0.001001 + 0.0005i$ . Computed results from quadratic Finite Element method and second order enrichment Mapped Infinite Partition of Unity Method are compared in figure 8.11. The distribution of the pressure is illustrated from the fan plane to the lip of the nacelle because the noise is completely attenuated by the liner.



**Fig. 8.11.** MIPUM 2-15-2 with 30828 degrees of freedom (top) and QFEM with 105016 degrees of freedom (bottom): Radiation of mode (26, 1) at  $ka = 30$  with admittance: (a) real part of the pressure (b) absolute value of pressure. The distribution of the pressure is illustrated from the fan plane to the lip of the nacelle.

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